

GRADIENT ESTIMATES AND GRADIENT BLOW UP  
OF VISCOSITY SOLUTIONS TO DEGENERATE ELLIPTIC  
EQUATIONS ON “INTERIOR BOUNDARIES”

Georgi Chobanov, Nikolay Kutev

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**Abstract**

Sharp conditions for gradient estimates or gradient blow up of the viscosity solutions to linear degenerate elliptic equations on the “interior boundaries” in bounded domains in  $\mathbb{R}^n$  are given.

**Key words:** Linear degenerate elliptic equations, viscosity solutions, gradient estimates, gradient blow up

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**1. Introduction.** Let

$$(1) \quad Lu = - \sum_{i,j=1}^n a^{ij}(x)u_{x_i x_j} + \sum_{i=1}^n b^i(x)u_{x_i} + c(x)u = f(x) \quad \text{in } \Omega$$

be a second order partial differential equation with non-negative characteristic form (degenerate elliptic equation)

$$(2) \quad \sum_{i,j=1}^n a^{ij} \xi_i \xi_j \geq 0 \quad \text{for } x \in \bar{\Omega}, \xi \in \mathbb{R}^n \setminus \{0\}$$

defined in a bounded region  $\Omega \subset \mathbb{R}^n$ ,  $n \geq 2$ . In this paper the regularity of viscosity solutions to such equations in a neighbourhood of the so-called “interior boundary” is studied. As it is well known, see paper [1], the correct way to give boundary conditions is determined by means of the so-called Fichera function

$$(3) \quad \beta(x) = \sum_{k=1}^n \left( b^k(x) + \sum_{j=1}^n a_{x_j}^{kj}(x) \right) \nu_k \quad \text{on } \Gamma,$$

where  $\Gamma$  is the characteristic part of the boundary. Also of interest, however, is the case when  $\Gamma$  is a smooth surface interior to  $\Omega$  dividing it in two subdomains  $\Omega_1$  and  $\Omega_2$  so that  $\Omega = \Omega_1 \cup \Gamma \cup \Omega_2$ ,  $\Gamma$  is characteristic for (1) and, moreover,  $\beta(x) = 0$  on  $\Gamma$ . In what follows we assume

$$(4) \quad a^{ij}, b^i, c, f, \psi, \partial\Omega, \Gamma \in C^\infty,$$

$$(5) \quad c(x) \geq c_0 > 0 \quad \text{in} \quad \bar{\Omega}, c_0 = \text{const.}$$

Moreover, if  $\nu(x)$  is the unit interior normal to  $\Gamma$  with respect to  $\Omega_1$ , it is assumed that degeneration occurs in the normal directions  $\nu(x)$  only, i.e.

$$(6) \quad \sum_{i,j=1}^n a^{ij}(x) \xi_i \xi_j = 0 \quad \text{iff} \quad x \in \Gamma \quad \text{and} \quad \xi = \nu(x).$$

In a previous paper [2] we proved that under conditions (2),  $\beta(x) = 0$  on  $\Gamma$  and (4)–(6), if the generalized curvature  $\mathcal{H}_\Gamma$  (introduced by J. SERRIN, see ch. II, sec. 9 of [3] or (13) below) of  $\Gamma$  has constant sign, i.e.  $\mathcal{H}_\Gamma \geq 0$  or  $\mathcal{H}_\Gamma \leq 0$  for every  $x \in \Gamma$ , then the following problems:

$$(7) \quad Lu_1 = f \quad \text{in} \quad \Omega_1,$$

$$(8) \quad Lu_2 = f \quad \text{in} \quad \Omega_2, \quad u = \psi \quad \text{on} \quad \partial\Omega$$

have unique viscosity solutions  $u_1 \in C(\bar{\Omega}_1)$ ,  $u_2 \in C(\bar{\Omega}_2)$ . Moreover,  $u_1 = u_2$  on  $\Gamma$  and the function  $U(x) = u_1(x)$  in  $\bar{\Omega}_1$  and  $U(x) = u_2(x)$  in  $\bar{\Omega}_2$  is the unique viscosity solution of equation (1) and

$$(9) \quad u = \psi \quad \text{on} \quad \partial\Omega.$$

Assuming that the equation degenerates on  $\Gamma$ , only from the hypoellipticity results in Th. 5.9 in [4] (see also [5, 6]), it follows that  $U(x) \in C^\infty(\Omega \setminus \Gamma)$ . In [2] we proved that  $U(x)$  is Hölder continuous on  $\Gamma$  if  $\mathcal{H}_\Gamma = 0$ .

The aim of the present paper is to prove Lipschitz regularity or gradient blow up of  $u_1$ ,  $u_2$  or equivalently  $U$  on  $\Gamma$  under sharp conditions. The results are obtained under weaker assumptions than the usual [7].

One motivation to study this problem comes from the optics in the reconstruction of a wave in Fresnel (paraxial) regime from its intensity (see [8] and the references there). The corresponding equation is

$$(10) \quad -\nabla (A^2(x, z) \nabla u) = \frac{\partial A^2(x, z)}{\partial z} \quad \text{in} \quad \Omega \subset R^2,$$

where  $A(x, z)$  is the amplitude of the wave and  $u(x)$  is the unknown phase of the wave. Since  $A = 0$  on  $\partial\Omega$ , easy calculations give us that conditions (2),  $\beta(x) = 0$  on  $\Gamma$  and (4)–(6) are satisfied in  $\Omega$  and on  $\partial\Omega$ .

The same type of degeneracy on some part of the boundary can be observed for the equation of BLACK–SCHOLES [9].

Another motivation is the so-called “interior gradient blow up” for the viscosity solutions of fully nonlinear, uniformly elliptic equations, this phenomenon being investigated in Theorem 2.9 and Corollary 2.12 in [10].

The paper is organized as follows: in Section 2 some definitions of viscosity solutions of equation (1) and of boundary value problem (1), (9) are given. Section 3 deals with gradient estimates of viscosity solutions on the “interior boundary”  $\Gamma$  while section 4 is dedicated to gradient blow up of the viscosity solutions.

**2. Preliminaries.** Throughout this paper we use the notations and definitions from [11] for equation (1) with the notation

$$(11) \quad F(x, u, p, X) = -\text{trace}(A(x)X) + \langle B(x), p \rangle + c(x)u,$$

where  $A(x) = \{a_{ij}(x)\}$ ,  $B(x) = (b_i(x))$ . Note that (5) is now equivalent to

$$(12) \quad F(x, r, p, X) - F(x, s, p, X) \geq c_0(r - s) \quad \text{for } r \geq s,$$

in  $\bar{\Omega}$ ,  $c_0 = \text{const} > 0$ .

We rely on the definitions of viscosity sub- and supersolutions from [11], Definition 2.2.

We say that the Dirichlet condition is satisfied in the classical sense, if  $u \in C(\bar{\Omega})$  and if  $u(x) = \psi(x)$  for every  $x \in \partial\Omega$ . Unfortunately the classical Dirichlet problem is not stable under small perturbations of the differential equation, see Section 7 in [11]. Therefore one has to weaken condition (9) so it can encompass small perturbations, too. To this end we use Definition 7.4 in [11].

Further on we will use the following comparison principle, Th 3.2 in [2] (see also Ex. 3.6 in [11]).

**Proposition 1** (Comparison principle). *Suppose (2),  $\beta(x) = 0$  on  $\Gamma$  and (4)–(6) hold. If  $u \in USC(\bar{\Omega})$   $v \in LSC(\bar{\Omega})$  are viscosity sub- and supersolution of (1) and  $u \leq v$  on  $\partial\Omega$  then  $u \leq v$  in  $\bar{\Omega}$ .*

Let us recall the definition of the generalized mean curvature of some  $C^2$  smooth surfaces introduced by J. Serrin in ch. II., sect. 9 of [3]. If  $\lambda^\tau(x)$ ,  $\kappa_\tau(x)$  are the principal directions and the principal curvatures of  $\Gamma$  at some point  $x \in \Gamma$  and  $\nu(x)$  is the interior unit normal to  $\Gamma$  (with respect to  $\Omega_1$ ), then

$$(13) \quad \mathcal{H}(x) = \lambda^\tau A \lambda^\tau \kappa_\tau + \nu A \nu H$$

is the generalized mean curvature of  $\Gamma$  at the point  $x$ . Here  $A = \{a^{ij}(x)\}$  and  $H = (\kappa_1 + \dots + \kappa_{n-1})/(n-1)$  is the ordinary mean curvature of  $\Gamma$ .

From (6) we have  $\nu A\nu = 0$  on  $\Gamma$  so that for equation (1)

$$(14) \quad \mathcal{H}(x) = \lambda^\tau A \lambda^\tau \kappa_\tau.$$

By means of the Perron method, the following existence and uniqueness result was proved in [2].

**Proposition 2** (Th. 4.1 in [2]). *Suppose (2),  $\beta(x) = 0$  on  $\Gamma$  and (4)–(6) hold and  $\mathcal{H}(x) \leq 0$  or  $\mathcal{H}(x) \geq 0$  for every  $x \in \Gamma$ . Then (7) has a unique viscosity solution  $u_1(x) \in C(\overline{\Omega}_1)$  which satisfies equation (7) in  $\overline{\Omega}_1$  without any boundary condition on  $\partial\Omega_1$ .*

Analogously, (8) has a unique viscosity solution  $u_2(x) \in C(\overline{\Omega}_2)$  satisfying equation (8) in  $\overline{\Omega}_2 \setminus \partial\Omega$  and the Dirichlet condition on  $\partial\Omega$  in classical sense.

Finally,  $u_1(x) = u_2(x)$  on  $\Gamma$  and the function  $U(x) = u_1(x)$  in  $\Omega_1$ ,  $U(x) = u_2(x)$  in  $\overline{\Omega}_2$  is the unique viscosity solution of (1), (9).

It is proved in Section 5 in [2] that the common value  $u_0(x)$  of  $u_1(x)$  and  $u_2(x)$  on  $\Gamma$ , which plays an important role of the considerations below, is the solution of the uniformly elliptic equation

$$(15) \quad L_0 u = - \sum_{\sigma, \tau=1}^{n-1} A^{\sigma\tau}(x) u_{\lambda_\tau \lambda_\sigma} + \sum_{\tau=1}^{n-1} B^\tau(x) u_{\lambda_\tau} + c(x)u - f(x) = 0$$

for  $x \in \Gamma$ . Equation (15) is obtained using a principal coordinate system at  $x_0 \in \Gamma$  consisting of the unit normal  $\nu$  and the principle directions  $\lambda^\tau \in T_\Gamma(x_0)$  ( $\tau = 1, \dots, n-1$ ). Taking into account (6), we get also

$$(16) \quad a^{ij}(x) \nu_i \nu_j = 0, \quad a^{ik}(x) \nu_i = 0, \quad a_{x_s}^{ij}(x) \nu_i \nu_j = 0$$

for  $k, s = 1, \dots, n$  and every  $x \in \Gamma$  because the function  $a^{ij}(x) \xi_i \xi_j$  has a minimum in  $\overline{\Omega} \times \mathbb{R}^n$  at the points  $x \in \Gamma$ ,  $\xi = \nu(x)$ .

Problem (15) has a unique classical solution  $u_0(x)$  on  $C^\infty$  manifold  $\Gamma$  under conditions (5)–(6).

The following statement is a revised version of Th. 5.1 in [2] treating Hölder continuity of  $u_1(x)$  and  $u_2(x)$  on  $\Gamma$ . The difference is that the condition  $\mathcal{H}_\Gamma = 0$  is replaced with the weaker requirement  $\mathcal{H}_\Gamma$  not to change its sign on  $\Gamma$ .

**Proposition 3.** *Suppose (2),  $\beta(x) = 0$  on  $\Gamma$  and (4)–(6) hold and  $\mathcal{H}_\Gamma \geq 0$  or  $\mathcal{H}_\Gamma \leq 0$  for every  $x \in \Gamma$ . Then the viscosity solutions  $u_1(x) \in C(\overline{\Omega}_1)$  and  $u_2(x) \in C(\overline{\Omega}_2)$  of problems (7) and (8) defined in Prop. 2.2 satisfy the boundary data  $u_0(x)$  on  $\Gamma$ , i.e.  $u_1(x) = u_2(x) = u_0(x)$  on  $\Gamma$ . Moreover, the viscosity solution  $U(x) = u_1(x)$  in  $\Omega_1$ ,  $U(x) = u_2(x)$  in  $\overline{\Omega}_2$  of (1), (9) is Hölder continuous on  $\Gamma$  with exponent  $\lambda \in (0, 1)$  depending on  $\|a^{ij}\|_{C^2(\overline{\Omega})}$ ,  $\|b^i\|_{C^1(\overline{\Omega})}$ ,  $c_0$  and  $\Gamma$ .*

**Proof.** Without loss of generality suppose that  $\mathcal{H}_\Gamma \geq 0$ . Let  $G_i = \Omega_i \cap \{d(x) < d_0\}$ ,  $d(x) = \text{dist}(x, \Gamma)$ , where  $d_0$  is sufficiently small so that  $d(x) \in C^2(G_i)$ ,  $i = 1, 2$ . We consider in  $G_1$  the barrier function  $v(x) = u_0(x) + Nd^\lambda(x)$ ,

where  $N = [\sup_{\Gamma} |u_0(x)| + \sup_{\overline{\Omega}} |u(x)|] d_0^{-\lambda}$  and  $u_0(x)$  is extended as a constant along the normals to  $\Gamma$ . If the positive constant  $\lambda$  is chosen sufficiently small, then  $v(x)$  is a classical supersolution of (7). Easy calculations in a principle coordinate system give

$$L_1 v \geq -K_2 d + N[-\lambda K_3 + c_0] d^\lambda \geq 0$$

in  $\overline{G}_1$  when  $\lambda$  is sufficiently small. Here  $f_1(x) = f(x)$  on  $\Gamma$  and is extended as a constant along the normal to  $\Gamma$ ,  $f_1 \in C^1(\overline{G}_1)$ ,  $K_2 = \sup_{\overline{G}_1} |\nabla(f - f_1)|$ ,  $K_3 = \sup_{\overline{G}_1} |\nabla(b^i \nu_i + \lambda^\tau A \lambda^\tau k_\tau)|$ ,  $A = \{a^{ij}\}$  and  $\lambda^\tau$ ,  $k_\tau$  are principle directions and principal curvatures of  $\Gamma$  at the point  $y(x) \in \Gamma$ , nearest to  $x \in G_1$ .

Since  $\mathcal{H}_\Gamma \geq 0$  from Lemma 4.2 and Remark 4.3 in [2], it follows that  $v(x)$  is a viscosity supersolution of (7) in  $G_1 \cup \Gamma$ , such that  $v \geq u_1$  on  $\partial G_1 \cap \{d(x) = d_0\}$  and the comparison principle, Lemma 3.4 in [2], yields  $v(x) \geq u_1(x)$  in  $\overline{G}_1$ . In the same way one can prove that  $w(x) = u_0(x) - Nd^\lambda(x)$  is a viscosity subsolution of (7) and hence  $u_2(x) \geq u_0(x) - Nd^\lambda(x)$  in  $\overline{G}_2$ , i.e.  $u_0(x) \leq u_2(x) \leq u_1(x) \leq u_0(x)$  for  $x \in \Gamma$  and hence  $u_1(x) = u_2(x) = u_0(x)$  in  $\Gamma$ . Now by means of the standard comparison principle, Prop. 2.3, one can prove  $u_1(x) \geq u_0(x) - Nd^\lambda(x)$  in  $\overline{G}_1$ ,  $u_2(x) \leq u_0(x) + Nd^\lambda(x)$  in  $\overline{G}_2$  without curvature conditions for  $\Gamma$ , which gives Prop. 3.

**3. Gradient estimates of the viscosity solutions on the “interior boundary”.** As Prop. 3 shows, the viscosity solutions  $u_1(x)$ , resp.  $u_2(x)$ ,  $U(x)$  are only Hölder continuous functions in  $\overline{\Omega}_1$ , resp. in  $\overline{\Omega}_2, \overline{\Omega}$ . In  $\Omega \setminus \Gamma$  these solutions are  $C^\infty$  smooth (see Th. 5.9 in [4] or [5]) and the open question is their higher regularity in a neighbourhood of  $\Gamma$ .

Let us mention that a sufficient condition for gradient estimates of the generalized  $L^p(\Omega)$  solutions of (1), (9) is proved in ([7], Th. 1.8.1) by means of elliptic regularization. The main condition is  $\sup_{x \in E_\delta} B_1(x) > 0$ , where  $E_\delta$  is a neighbourhood of the set  $E$  of degeneracy of  $L$ , i.e.  $E = \{x \in \Omega; \det\{a^{ij}(x)\} = 0\}$ ,

$$B_1(x) = c(x) - \frac{1}{4} M n + \max_s \left\{ b_{x_s}^s - \frac{1}{2} \sum_{j=1}^n |b_{x_s}^j| - \frac{1}{2} \sum_{j=1, j \neq s}^n |b_{x_j}^s| \right\}$$

and  $M = \min_{1 \leq s \leq n} M_s$  for

$$M_s = \inf_{v \in W^{2,\infty}(\Omega)} \inf_{x \in \overline{E}_\delta} \left( \sum_{k,j=1}^n a_{x_s}^{kj} v_{x_k x_j} \right)^2 / \left( \sum_{i,j,k,s=1}^n a^{kj} v_{x_k x_s} v_{x_j x_s} \right).$$

We will give precise conditions for gradient estimates or gradient blow up of the viscosity solutions on  $\Gamma$ .

**Theorem 1.** *Suppose (2),  $\beta(x) = 0$  on  $\Gamma$  and (4)–(6) hold and  $\mathcal{H}_\Gamma \geq 0$  or  $\mathcal{H}_\Gamma \leq 0$  for every  $x \in \Gamma$ . If*

$$(17) \quad \frac{\partial b^i}{\partial \nu}(x) \nu_i + \lambda^\tau \frac{\partial A}{\partial \nu}(x) \lambda^\tau \kappa_\tau + \lambda^\tau A(x) \lambda^\tau \kappa_\tau^2 + c(x) \geq c_1,$$

$c_1 = \text{const} > 0$  for every  $x \in \Gamma$  then the viscosity solutions  $u_1(x) \in C(\overline{\Omega}_1)$  and  $u_2(x) \in C(\overline{\Omega}_2)$  of problems (7) and (8) defined in Prop. 2 are Lipschitz continuous in the neighbourhood of  $\Gamma$ .

SKETCH OF THE PROOF. Using the same technique as in obtaining (15) and (16), in  $G_i = \Omega_i \cap \{d(x) < d_0\}$  equation (1) can be rewritten in the form

$$(18) \quad Lu = -A^{nn}(x)u_{\nu\nu} - A^{\tau\sigma}(x)u_{\lambda^\tau\lambda^\sigma} - A^{\tau n}(x)u_{\lambda^\tau\nu} - A^{n\tau}(x)u_{\nu\lambda^\tau} \\ + B^n(x)u_\nu + B^\tau(x)u_{\lambda^\tau} + c(x)u - f(x) = 0.$$

Moreover, from (3), (16) and the derivatives of (16) in  $\lambda^\tau$  directions on  $\Gamma$ , we get

$$(19) \quad 0 = \beta(x) = b^i(x)\nu_i + \mathcal{H},$$

$$(20) \quad A^{nn}(x) = A^{\tau n}(x) = A^{n\tau}(x) = B^n(x) = 0$$

for every  $x \in \Gamma$  and  $\tau = 1, \dots, n-1$ . Let us consider in  $G_1$  the barrier function  $v(x) = u_0(x) + Nd(x)$ . Here the constant  $N$  will be chosen sufficiently large and the classical solution  $u_0(x)$  of (15) is extended as a constant along the normals to  $\Gamma$  in  $G_i$ . As in the proof of Prop 2.5. we get

$$(21) \quad L_1v \geq -K_1d + N [b^i(x)\nu_i + \lambda^\tau A(x)\lambda^\tau k_\tau + \lambda^\tau A(x)\lambda^\tau k_\tau^2 d + c(x)d - K_2d^2].$$

Here  $f_1(x) = f(x)$  on  $\Gamma$  and  $f_1(x)$  is extended as a constant along the normals to  $\Gamma$  in  $G_i$ ,  $f_1 \in C^1(\overline{G}_i)$ ,  $K_1 = \sup_{\overline{G}_i} |\nabla(f - f_1)|$ ,  $K_2 = \sup_{\overline{G}_i} \left| \frac{\lambda^\tau A(x)\lambda^\tau k_\tau^3}{1 - k_\tau d} \right|$ . When  $N$  is sufficiently large and  $d_0$  is sufficiently small, we get  $L_1v \geq 0$  in  $G_1$ ,  $v = u_1$  on  $\Gamma$ ,  $v \geq u_1$  for every  $x \in \partial G_1 \cap \{d(x) = d_0\}$  and from the comparison principle, Proposition 1, it follows that  $u_1(x) \leq v(x) = u_0(x) + Nd(x)$  for every  $x \in \overline{G}_1$ . In the same way one can prove that the function  $w(x) = u_0(x) - Nd(x)$  is a viscosity subsolution of (7) in  $G_1$ , i.e.

$$(22) \quad u_0(x) - Nd(x) \leq u_1(x) \leq u_0(x) + Nd(x) \quad \text{for every } x \in \overline{G}_1.$$

Inequalities (22) prove Theorem 1 for  $u_1(x)$ . The proof of the Lipschitz continuity of  $u_2(x)$  is similar and is omitted here.

**4. Gradient blow up of the viscosity solutions on the “interior boundary”.** In this section we will show that condition (17) in Theorem 3.1 is sharp.

**Theorem 2.** Suppose (2),  $\beta(x) = 0$  on  $\Gamma$  and (4)–(6) hold and  $\mathcal{H}_\Gamma \geq 0$  or  $\mathcal{H}_\Gamma \leq 0$  for every  $x \in \Gamma$ . If

$$(23) \quad \frac{\partial b^i}{\partial \nu}(x_0)\nu_i + \lambda^\tau \frac{\partial A}{\partial \nu}(x_0)\lambda^\tau \kappa^\tau + \lambda^\tau A(x_0)\lambda^\tau \kappa_\tau^2 + c(x_0) < 0$$

for some  $x_0 \in \Gamma$ ,

$$(24) \quad f(x) \equiv 0 \quad \text{on} \quad \Gamma, f(x) \not\equiv 0 \quad \text{in} \quad \Omega_1$$

and either

$$(25) \quad f(x) \leq 0 \quad \text{in} \quad \Omega_2 \quad \text{or} \quad f(x) \geq 0 \quad \text{in} \quad \Omega_2,$$

then the gradient of the viscosity solution  $u_1(x_0)$  blows up at  $x_0$  in direction  $\nu(x_0)$  where  $\nu(x_0)$  is the interior normal to  $\Gamma$  at  $x_0$  with respect to  $\Omega_1$ . More precisely

$$(26) \quad \lim_{t \rightarrow +0} \frac{u_1(x_0 + t\nu(x_0)) - u_1(x_0)}{t} = -\infty \quad \text{or} \quad \lim_{t \rightarrow +0} \frac{u_1(x_0 + t\nu(x_0)) - u_1(x_0)}{t} = \infty$$

respectively.

SKETCH OF THE PROOF. Since  $f(x) \equiv 0$  on  $\Gamma$  the unique solution of (15) is  $u_0 = 0$ . Suppose that (25) holds. Then the function  $\bar{u} = 0$  is a classical supersolution of (7) in  $\Omega_1$  because

$$L_1 \bar{u} = -f \geq 0 = L_1 u_1$$

in  $\Omega_1$  and  $\bar{u} = u_1$  on  $\Gamma$ . From the comparison principle, Prop. 1, it follows that  $u_1 \leq \bar{u} = 0$  in  $\bar{\Omega}_1$ . If  $u_1(y) = 0$  at some interior point  $y \in \Omega_1$ , then from the strong interior maximum principle it follows that  $u_1 \equiv 0$  on every compact  $K \subset \Omega_1$ ,  $y \in K$ , and hence  $f \equiv 0$  in  $K$  contrary to (24). Note that from the hypoellipticity of equation (1) in  $\Omega_1$  we have  $u_1 \in C^\infty(\Omega_1)$ . Hence

$$(27) \quad u_1(x) < 0 \quad \text{in} \quad \Omega_1, \quad u_1(x) = 0 \quad \text{on} \quad \Gamma.$$

From (23) there exist a constant  $c_2 > 0$  and a ball  $B_0(x_0, r_0) = \{|x - x_0| < r_0\}$  such that

$$\frac{\partial b^i}{\partial \nu}(x) \nu_i + \lambda^\tau A_\nu(x) \lambda^\tau \kappa_\tau^2 + \lambda^\tau \frac{\partial A}{\partial \nu}(x) \lambda^\tau \kappa_\tau + c(x) \leq -2c_2$$

for every  $x \in B_0(x_0, r_0) \cap \bar{\Omega}_1$ . Here the normal  $\nu$  and the principle directions  $\lambda^\tau$  are calculated at the point  $y(x) \in \Gamma$  nearest to  $x$ .

Let  $\Gamma_1$  be an internal touching surface to  $\partial\Omega_1$  at the point  $x \in \Gamma$ ,  $\Gamma_1 \subset \Omega_1$ ,  $\Gamma_1 \cap \Gamma = \{x\}$ . We consider a domain  $G \subset \Omega_1$  with boundary  $\partial G = \Gamma_1 \cup \Gamma_0$ ,  $\Gamma_0 \subset \partial B_0$ . If  $n$ ,  $\mu^\tau$ ,  $k_\tau$  are the unite interior normal to  $\Gamma_1$  (with respect to  $G$ ) and the principle directions and the principle curvatures of  $\Gamma_1$ , we suppose that  $\Gamma_1$  is sufficiently close to  $\Gamma$  and  $r_0$  is sufficiently small so that  $\lambda^\tau(x_0) = \mu^\tau(x_0)$ ,  $\kappa_\tau(x_0) = k_\tau(x_0)$  and the inequality

$$(28) \quad \frac{\partial b^i}{\partial n}(x) n_i + \mu^\tau A(x) \mu^\tau k_\tau^2 + \mu^\tau \frac{\partial A}{\partial n}(x) \mu^\tau k_\tau + c(x) \leq -c_2$$

holds for every  $x \in G$ . Here  $n$ ,  $\mu^\tau$ ,  $k_\tau$  are calculated at the point  $y(x) \in \Gamma_1$  nearest to  $x$ .

Let  $\rho(x) = \text{dist}(x, \Gamma_1)$ ,  $\Gamma_0 = \partial G \cap \partial B_0$  and let us define  $\sup_{x \in \Gamma_0} u_1(x) = -M$ ,  $K_1 = \sup_G \frac{nA(x)n}{\rho^2(x)}$ ,  $K_2 = \sup_G \left| \frac{\mu^\tau A(x) \mu^\tau k_\tau^3}{1 - k_\tau d} \right|$ . Note that  $M > 0$  from (27) while  $K_1 < \infty$  from (16) and the choice of  $\Gamma_1$ .

We consider the barrier function  $v = -\epsilon \rho^\lambda(x)$  in  $G$ , where  $\sup_G \rho(x) \leq \frac{C_2}{4K_2}$  is fixed,  $\epsilon$  is sufficiently small and  $\lambda$  is sufficiently close to 1.

Repeating the calculations in the proof of Proposition 3 we get

$$(29) \quad L_1 v \geq \epsilon \rho^\lambda \left[ -\frac{C_2}{4} + \frac{C_2}{2} - \frac{C_2}{4} \right] = 0.$$

Since  $v \geq -\epsilon \sup_{x \in \Gamma_1} \rho^\lambda(x) = -M \geq u_1$  for  $x \in \Gamma_1$  and  $v = 0 \geq u_1$  for  $x \in \Gamma_0$  from the comparison principle, Proposition 1, it follows that

$$(30) \quad -\epsilon \rho^\lambda(x) = v(x) \geq u_1(x) \quad \text{for every } x \in G.$$

Hence for every  $x = x_0 + \rho n(x_0) = x_0 + \rho(x) \nu(x_0)$  we get from (30)

$$\lim_{t \rightarrow +0} \frac{u_1(x_0 + t\nu(x_0)) - u_1(x_0)}{t} = \lim_{t \rightarrow +0} \frac{u_1(x_0 + t\nu(x_0))}{t} \leq -\epsilon \lim_{t \rightarrow +0} \frac{t^\lambda}{t} = -\infty$$

which proves (26).

## REFERENCES

- [1] FICHERA G. *Atti Accad. Naz. Lincei, Memorie (VIII)*, **5**, 1956, No 8, 1–30.
- [2] CHOBANOV G., N. KUTEV. *Mediterr. J. Math.*, **9**, 2012, No 4, 789–801, DOI: 10.1007/s00009-011-0151-7.
- [3] SERRIN J. R. *Soc. Lond. Philos. Trans. Ser. A Math. Phys. Eng. Sci.*, **264**, 1969, 413–496.
- [4] HÖRMANDER L. *Math. Ann.*, **217**, 1975, No 2, 165–188.
- [5] POPIVANOV P. *Rend. Sem. Mat. Univ. Politec. Torino*, **66**, 2008, No 4, 321–337.
- [6] RADKEVIČ E. *DAN SSSR*, **187**, 1969, 274–277; *Soviet Math. Dokl.*, **10**, 1969, 849–853 (English translation).
- [7] OLEINIK O. A., E. RADKEVICH. *Itogi Nauki*, Moscow, 1971.
- [8] LEE CHUNG-MIN, J. RUBINSTEIN. *Quart. Appl. Math.*, **64**, 2006, No 4, 735–747.
- [9] BLACK F., M. SCHOLES. *J. Political Economy*, **81**, 1973, 637–654.
- [10] KAWOHL B., N. KUTEV. *Acta Mathematica Scientia*, 32B(1), 2012, 15–40.
- [11] CRANDALL M., H. ISHII, P.-L. LIONS. *Bull. Amer. Math. Soc. (N.S.)*, **27**, 1992, No 1, 1–67.

*Institute for Mathematics and Informatics  
Bulgarian Academy of Sciences  
Acad. G. Bonchev Str., Bl. 8  
1113 Sofia, Bulgaria  
e-mail: chobanov@math.bas.bg*